On the use of fractional derivative operators to describe viscoelastic damping in structural dynamics - FE formulation of sandwich beams and approximation of fractional derivatives by using the  $G^{\alpha}$  scheme

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Dérivation fractionnaire en mécanique - Etat-de-l'art et applications

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# **Objectives**

• To implement a fractional derivative model into a finite element code in structural dynamics



 To test a finite difference based on scheme to approximate fractional derivative operators in viscoelastic constitutive equations



# Outline

I. Implementation of the fractional Zener model into a FE code – application to sandwich beams

- Constitutive equations fractional Zener model
- Grünwald-Letnikov approximation
- FE formulation & algorithm
- Examples: validation & results

#### II. Testing the $G^{\alpha}$ scheme

- Outline
- Elementary tests
- Error estimates
- The damped oscillator problem
- Viscoelastic cantilever beam

I. Implementation of the fractional Zener model into a FE code – application to sandwich beams

# Fractional derivative viscoelastic model One-dimensional constitutive equation

 $\Rightarrow$  Fractional derivative Zener Model (4 parameters)

R.L. Bagley and P.J. Torvik, AIAA (1983)

$$\begin{array}{ll} \mbox{In frequency domain} & E^*(\omega) = \frac{E_o + E_\infty (i\omega\tau)^\alpha}{1 + (i\omega\tau)^\alpha} & 0 < \alpha < 1 \\ & E_\infty > E_o \\ \mbox{In time domain} & \sigma(t) + \tau^\alpha D^\alpha \sigma(t) = E_o \varepsilon(t) + \tau^\alpha E_\infty D^\alpha \varepsilon(t) & \tau > 0 \\ & \mbox{with} & D^\alpha \sigma(t) = \frac{1}{\Gamma(1-\alpha)} \frac{\mathrm{d}}{\mathrm{d}t} \int_0^t \frac{\sigma(s)}{(t-s)^\alpha} \mathrm{d}s \end{array}$$

### Master curves of the ISD112 (3M<sup>™</sup>) at 27°

#### Fractional model versus classical Zener model



Identified model parameters:

 $E_o = 1.5$  MPa,  $E_{\infty} = 69.9$  MPa,  $\alpha = 0.7915$ ,  $\tau = 1.4052 \times 10^{-2}$  ms

### Grünwald approximation & time discretization

 $\sigma(t) + \tau^{\alpha} D^{\alpha} \sigma(t) = E_o \varepsilon(t) + \tau^{\alpha} E_{\infty} D^{\alpha} \varepsilon(t)$ 

Approximation for fractional derivativesJ. Padovan, Comp. Mech. (1987) $(D^{\alpha}f)_n \approx \frac{1}{\Delta t^{\alpha}} \sum_{j=0}^{N_t} A_{j+1} f_{n-j}$ A. Schmidt and L. Gaul, Proc. 2nd ECCMR (2001)with $A_{j+1} = \frac{\Gamma(j-\alpha)}{\Gamma(-\alpha)\Gamma(j+1)}$ Variable changingOnly one fractional derivative term

Discretizing...

$$\bar{\varepsilon}_i^{n+1} = (1-c)\frac{E_{\infty} - E_o}{E_{\infty}}\varepsilon_i^{n+1} - c\sum_{j=1}^{N_t} A_{j+1}\bar{\varepsilon}_i^{n+1-j} \qquad \text{with} \qquad c = \frac{\tau^{\alpha}}{\tau^{\alpha} + \Delta t^{\alpha}}$$

### Finite element formulation

- Sandwich beam element
- Fractional model implementation
- Dynamic time integration

The structure consists on a three-layer beam  $\Rightarrow$  constrained layer damping treatment



### Sandwich-beam kinematics

Mead and Markus, JSV (1969), Trindade, Benjeddou and Ohayon, IJNME (2001)



Constrained layer damping treatment

- Laminated elastic faces: Euler-Bernoulli beam
- Viscoelastic core: Timoshenko beam

### Viscoelastic core



"Anelastic displacements" are evaluated as functions of their own time histories

$$\bar{\mathbf{q}}_{e}^{n+1} = (1-c) \frac{E_{\infty} - E_{o}}{E_{\infty}} \mathbf{q}_{e}^{n+1} - c \sum_{j=1}^{N_{t}} A_{j+1} \bar{\mathbf{q}}_{e}^{n+1-j}$$
  
Elementary dof vector  $\mathbf{q}^{e} = [\bar{u}_{1} \ w_{1} \ w'_{1} \ \tilde{u}_{1} | \bar{u}_{2} \ w_{2} \ w'_{2} \ \tilde{u}_{2} ]^{\mathrm{T}}$ 

### **Equation of motion**

$$\mathbf{M}_e \ddot{\mathbf{q}}_e^{n+1} + (\mathbf{K}_e + \bar{\mathbf{K}}_e)\mathbf{q}_e^{n+1} = \mathbf{F}_e^{n+1} + \bar{\mathbf{F}}_e^{n+1}$$

+ initial conditions

#### Remarks

**D** For an elastic material ( $\tau = 0$ ), classical elastic stiffness matrix

is obtained

 $\bar{\mathbf{K}}_{c}^{e} = c \frac{E_{o} - E_{o}}{E_{o}} \mathbf{K}_{c}^{e} \qquad \qquad c = 0$ 

□ For the Zener model ( $\alpha$  = 1), previous equations correspond to a backward Euler discretization of the constitutive equations

$$c = \frac{\tau}{\tau^* + \Delta t}$$

### Dynamic time integration

 $\mathbf{M}\ddot{\mathbf{q}}_{n+1} + \mathbf{K}^*\mathbf{q}_{n+1} = \mathbf{F}_{n+1} + \bar{\mathbf{F}}_{n+1}$ 



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### Validation – Example 1

Viscoelastic fixed-free bar subjected to unit step load

M. Enelund et B.L. Josefson, AIAA Journal (1997)



### Validation – Example 2

Viscoelastic simply-supported Timoshenko beam subjected to a uniform step function load

T.M. Chen, IJNME (1995)



Validation for  $\alpha = 1$ 

good agreement with the "hybrid Laplace transform/FE method"

#### Sandwich beam subjected to a triangular pulse



#### Influence of the time step and truncation

Calculations performed with whole history



<u>Remark</u>: Energy balance is strictly zero.

#### Influence of the time step and truncation



 $D_{ref}$  Very accurate solution with a fine time discretization:  $N_{step} = 2500$   $N_t = 2500$ 

Existence of a truncation time where the recent history must be stored.

1000

1000

## II. Testing the $G^{\alpha}$ scheme

### How to obtain a fractional derivative operator

Euler backward operatorGear operator
$$E = \frac{1}{\Delta t} (I - \delta^-)$$
 $G = \frac{1}{\Delta t} \left[ \frac{3}{2}I - 2\delta^- + \frac{1}{2}(\delta^-)^2 \right]$ Grünwald-Letnikov operator $G^{\alpha}$  operator $GL = E^{\alpha} = \frac{1}{\Delta t^{\alpha}} (I - \delta^-)^{\alpha}$  $G^{\alpha} = \frac{1}{\Delta t^{\alpha}} \left( \frac{3}{2} \right)^{\alpha} \left[ I - \frac{4}{3}\delta^- + \frac{1}{3}(\delta^-)^2 \right]^{\alpha}$ 

### The $G^{\alpha}$ scheme

First five coefficients  $g_{j+1}$ 

The  $\alpha$ -derivative of u at time  $t^n$  can be approximated by:

$$(\mathbf{G}^{\alpha}u)^{n} = \frac{1}{\Delta t^{\alpha}} \left(\frac{3}{2}\right)^{\alpha} \sum_{j=0}^{\infty} \mathbf{g_{j+1}}u^{n-j}$$

| j | $\alpha = 1/3$         | $\alpha = 1/2$    | $\alpha = 3/4$   |
|---|------------------------|-------------------|------------------|
| 0 | 1                      | 1                 | 1                |
| 0 | 1                      | 1                 | 1                |
| 1 | $-\frac{4}{9}$         | $-\frac{2}{3}$    | -1               |
| 2 | $-\frac{7}{81}$        | $-\frac{1}{18}$   | $\frac{1}{12}$   |
| 3 | $-\frac{104}{2187}$    | $-\frac{1}{27}$   | $-\frac{1}{108}$ |
| 4 | $-\frac{643}{19683}$   | $-\frac{17}{648}$ | $-\frac{1}{96}$  |
| 5 | $-\frac{4348}{177147}$ | $-\frac{19}{972}$ | $-\frac{7}{864}$ |

### **Elementary tests**

Consider the following power function

$$u(t) = t^{\nu}, \quad \nu \ge 0$$

We search to calculate its  $\alpha$ -derivative by using an exact solution and an approximated one with the G<sup> $\alpha$ </sup>-scheme, such that:

Exact solution

$$\mathcal{D}^{\alpha}u(t) = \frac{\Gamma(\nu+1)}{\Gamma(\nu+1-\alpha)}t^{\nu-\alpha}$$

Approximated solution

$$(G^{\alpha}u)^{n} = \frac{1}{\Delta t^{\alpha}} \left(\frac{3}{2}\right)^{\alpha} \sum_{j=0}^{\infty} g_{j+1}u^{n-j}$$

How accurate is the  $G^{\alpha}$ -scheme for this case?

## Error estimates in $L^{\infty}$ norm



|                             | $\alpha = 1/3$    |      | $\alpha = 1/2$    |       | $\alpha = 3/4$    |       |
|-----------------------------|-------------------|------|-------------------|-------|-------------------|-------|
| Power                       | $G^{lpha}$        | GL   | $G^{lpha}$        | GL    | $G^{lpha}$        | GL    |
| 1/4                         | 0.25              | 0.25 | 0.25              | 0.25  | 0.25              | 0.25  |
| 1/3                         | 0.33              | 0.33 | 0.33              | 0.33  | 0.33              | 0.33  |
| 1/2                         | 0.50              | 0.50 | 0.50              | 0.50  | 0.50              | 0.50  |
| 2/3                         | 0.67              | 0.67 | 0.67              | 0.67  | 0.67              | 0.67  |
| 3/4                         | 0.75              | 0.75 | 0.75              | 0.75  | 0.75              | 0.75  |
| 1                           | 1.00              | 1.00 | 1.00              | 1.00  | 1.00              | 1.00  |
| 5/4                         | 1.25              | 1.25 | 1.25              | 1.25  | 1.25              | 1.25  |
| 4/3                         | 1.33              | 1.33 | 1.33              | 1.33_ | 1.33              | 1.33  |
| 3/2                         | 1.50              | 1.33 | 1.50              | 1.50  | 1.50              | 1.50  |
| 5/3                         | 1.67              | 1.33 | 1.67              | 1.50  | 1.67              | 1.67_ |
| 7/4                         | 1.75              | 1.33 | 1.75              | 1.50  | 1.75              | 1.75  |
| 2                           | 2.00              | 1.33 | 2.00              | 1.49  | 2.00              | 1.74  |
| 9/4                         | 2.25              | 1.32 | 2.25              | 1.49  | 2.25              | 1.74  |
| 7/3                         | 2.33 <sup>A</sup> | 1.32 | 2.33              | 1.49  | 2.33              | 1.74  |
| 5/2                         | 2.32              | 1.32 | 2.50 <sup>B</sup> | 1.48  | 2.50              | 1.73  |
| 8/3                         | 2.31              | 1.32 | 2.49              | 1.48  | 2.67              | 1.73  |
| 11/4                        | 2.30              | 1.32 | 2.48              | 1.48  | 2.75 <sup>C</sup> | 1.73  |
| 3                           | 2.29              | 1.31 | 2.46              | 1.48  | 2.73              | 1.72  |
| 13/4                        | 2.27              | 1.31 | 2.45              | 1.47  | 2.72              | 1.72  |
| 10/3                        | 2.27              | 1.31 | 2.45              | 1.47  | 2.71              | 1.72  |
| 7/2                         | 2.26              | 1.31 | 2.44              | 1.47  | 2.70              | 1.71  |
| 11/3                        | 2.25              | 1.30 | 2.43              | 1.47  | 2.69              | 1.71  |
| 15/4                        | 2.25              | 1.30 | 2.42              | 1.47  | 2.69              | 1.71  |
| 4                           | 2.24              | 1.30 | 2.41              | 1.46  | 2.67              | 1.71  |
| $G^{\alpha}$ : $\alpha + 2$ |                   |      | $GL: \alpha + 1$  |       |                   |       |

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### The damped oscillator problem

The problem is given such that

$$\begin{cases} m\ddot{u} + c\tau^{\alpha}\mathcal{D}^{\alpha}u + ku = f, \quad t > 0\\ u(0) = \dot{u}(0) = 0 \end{cases}$$



### Algorithm

We search  $u^{n+1}$  such that

$$m\ddot{u}^{n+1} + (k+\kappa)u^{n+1} = f^{n+1} + \phi^{n+1}$$

where 
$$\kappa = \frac{c\tau^{\alpha}}{\Delta t^{\alpha}} \left(\frac{3}{2}\right)^{\alpha}$$
 and  $\phi^{n+1} = -\frac{c\tau^{\alpha}}{\Delta t^{\alpha}} \left(\frac{3}{2}\right)^{\alpha} \sum_{k=1}^{N} g_{k+1} u^{n+1-k}$ 

Then we use the average acceleration algorithm:

• Prediction 
$$u_p = u^n + \Delta t \dot{u}^n + \frac{\Delta t^2}{4} \ddot{u}^n$$
  $\dot{u}_p = \dot{u}^n + \frac{\Delta t}{2} \ddot{u}^n$ 

• Evaluation 
$$\ddot{u}^{n+1} = \frac{1}{s} \left[ f^{n+1} + \phi^n - (k+\kappa)u_p^{n+1} \right]$$

• Correction 
$$u^{n+1} = u_p + \frac{\Delta t^2}{4}\ddot{u}^{n+1} \quad \dot{u}^{n+1} = \dot{u}_p + \frac{\Delta t}{2}\ddot{u}^{n+1}$$

• Calculation of 
$$\phi^{n+1}$$







The rate of convergence computed with the  $L^\infty$  norm is:

0.99 when using the GL-scheme

**1.90** when using the  $G^{\alpha}$ -scheme

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## Extension to viscoelastic beams

Constitutive equations

Bagley & Torvik, 1984



The 1D fractional Zener model is given by:

$$\sigma(t) + \tau^{\alpha} \mathcal{D}^{\alpha} \sigma(t) = E_o \varepsilon(t) + E_{\infty} \tau^{\alpha} \mathcal{D}^{\alpha} \varepsilon(t)$$

Introducing the strain-type internal variable  $\varepsilon_{\alpha} = \varepsilon - \frac{\sigma}{E_{\infty}}$ ,

We obtain only one fractional derivative operator:

$$\varepsilon_{\alpha}(t) + \tau^{\alpha} \mathcal{D}^{\alpha} \varepsilon_{\alpha}(t) = \frac{E_{\infty} - E_{o}}{E_{\infty}} \varepsilon(t)$$

Implementation in a finite element code and resolution of the dynamic problem

### Viscoelastic cantilever beam



We search to solve the problem

$$\left\{ egin{aligned} \mathbf{M}\ddot{\mathbf{q}}^{n+1} + (\mathbf{K}+\mathcal{K})\mathbf{q}^{n+1} &= \mathbf{F}^{n+1} + \mathbf{\Phi}^{n+1} \ \mathbf{q}^0 &= \dot{\mathbf{q}}^0 = \mathbf{0} \end{aligned} 
ight.$$

where 
$$\mathcal{K} = c_{\alpha} \frac{E_{\infty} - E_o}{E_o} \mathbf{K}$$
 and  $\Phi^{n+1} = -c_{\alpha} \frac{E_{\infty}}{E_o} \mathbf{K} \sum_{k=1}^{N} g_{k+1} \mathbf{q}_{\alpha}^{n+1-k}$   
with  $c_{\alpha} = \frac{\tau^{\alpha}}{\tau^{\alpha} + \Delta t^{\alpha}}$ 

### Algorithm

We search  $u^{n+1}$  such that

$$\begin{cases} \mathbf{M}\ddot{\mathbf{q}}^{n+1} + (\mathbf{K} + \mathcal{K})\mathbf{q}^{n+1} = \mathbf{F}^{n+1} + \Phi^{n+1} \\ \mathbf{q}^0 = \dot{\mathbf{q}}^0 = \mathbf{0} \end{cases}$$

Then we use the *average acceleration algorithm*, which requires the calculation:

• 
$$\Phi^{n+1} = -c_{\alpha} \frac{E_{\infty}}{E_o} \mathbf{K} \sum_{k=1}^{N} g_{k+1} \mathbf{q}_{\alpha}^{n+1-k}$$
  
• 
$$\mathbf{q}_{\alpha}^{n+1} = (1-c_{\alpha}) \frac{E_{\infty} - E_o}{E_{\infty}} \mathbf{q}^{n+1} - c_{\alpha} \sum_{k=1}^{N} g_{k+1} \mathbf{q}_{\alpha}^{n+1-k}$$

#### Finite element implementation of a viscoelastic beam submitted to a mechanical load



The whole history of the anelastic displacements is used in the calculations.

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# Quelques références

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## **Conclusions and Perspectives**

- ✓ Sandwich beam with viscoelastic core and elastic faces
- ✓ FE implementation of a viscoelastic fractional derivative model
- ✓ The G<sup> $\alpha$ </sup> scheme is no more costly than the Grünwald-Letnikov one and it appears to be ( $\alpha$ +2)-order accurate for power functions
- $\checkmark$  The combination Newmark-  $G^{\alpha}$  yields a two-order accuracy for the damped oscillator problem

Coming work: to investigate the order of convergence obtained by using the  $G^{\alpha}$  scheme in the case of sandwich beams